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TEMPERATURE DEPENDENCE OF ANHARMONIC EXAFS OSCILLATION OF CRYSTALLINE SILICON

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Article info	Abstract:
	In this work, the anharmonic extended X-ray absorption fine structure
Recieved:	(EXAFS) oscillation of crystalline silicon (c-Si) is presented in terms of the
16/9/2020	Debye-Waller factor using the cumulant expansion approach up to the fourth-
Accepted:	classical anharmonic correlated Einstein (ACE) model and suitable analysis
10/12/2020	procedure, in which thermodynamic parameters are derived from the
	anharmonic effective potential obtained using the first shell near-neighbor contribution approach. The analysis of the temperature dependence of the
Keywords:	EXAFS oscillation is performed via evaluating the influence of the cumulants on the amplitude reduction and the phase shift of the anharmonic EXAFS
EXAFS analysis; Debye-	oscillation. The numerical results are found to be in good agreement with those
Waller factor; Anharmonic correlated Einstein model; Crystalline silicon.	obtained using the quantum ACE model and experiments at various
	temperatures. The obtained results are useful in analyzing the experimental
	EXAFS data of c-Si.

1. Introduction

The extended X-ray absorption fine structure (EXAFS) has been developed into a powerful technique and is widely used to determine many structural parameters and dynamic properties of materials [1]. However, the position of atoms is not stationary, and their interatomic distance always changes due to thermal vibrations that were detected by Beni & Platzman [2]. They cause anharmonic effects on crystal vibrations and smear out the EXAFS oscillations. The anharmonicity of the potential yields additional terms in the EXAFS oscillation, so if ignoring these terms, can lead to non-negligible errors in the structural parameters. The use of the moments of the radial distribution function (or cumulants) to investigate local disorder of EXAFS spectra was introduced by Rehr [3] who showed that the Debye-Waller (DW) factor of EXAFS spectra has a natural cumulant expansion in powers of the photoelectron wavenumber. The connection between the DW factor and the EXAFS cumulants was described in detail in the cumulant expansion approach (*ratio* method) by Bunker [4]. The ratio method is particularly appealing because it summarizes the relevant structural and dynamic

information that is easily obtained from the experimental EXAFS spectra.

Recently, a classical anharmonic correlated Einstein (ACE) model [5] was developed based on the anharmonic effective (AE) potential [6] and the classical statistical theory [7]. This model has the advantage that the expressions of the first four EXAFS cumulants are obtained in explicit and simple forms, so it is very convenient for anharmonic EXAFS data analysis in the range of temperatures not too low. It has also been successfully applied to investigate the anharmonic EXAFS oscillation for diamond crystals by Tien et al. [8]. However, the temperature dependence of anharmonic EXAFS amplitude and phase has not been discussed in detail. Besides, it has been applied to successfully investigate crystalline germanium but has not yet been investigated crystalline silicon (c-Si). Therefore, the analysis of the temperature dependence of anharmonic EXAFS oscillation for c-Si will be a necessary addition to evaluate the effectiveness of the classical ACE model in the EXAFS technique.

2. Basic formulae of EXAFS function and anharmonic effective potential of c-Si

2.1. Basic formulae of EXAFS function

The EXAFS oscillation for a single coordination shell, including thermal disorders has the form:

$$\chi(k,T) = A(k,T)\sin\Phi(k,T), \qquad (1)$$

where A(k,T) and $\Phi(k,T)$ are the EXAFS amplitude and phase, respectively, and k is the photoelectron wavenumber, and T is the temperature.

The *K*-edge EXAFS oscillation for the distribution of identical atoms is described within the framework of single-scattering and plane-wave approximations. Following the approach proposed by Tien [9], the logarithm of amplitude ratio $M(k,T_1,T_2) = \ln[A(k,T_2)/A(k,T_1)]$ and the linear phase difference $\Delta\Phi(k,T_1,T_2) = \Phi(k,T_2) - \Phi(k,T_1)$ between temperatures T_2 and T_1 in the cumulant expansion approach up to the fourth-order, which are given as follows:

$$M(k,T_{1},T_{2}) \Box -2k^{2} \left\{ \sigma^{2}(T_{2}) - \sigma^{2}(T_{1}) \right\} + \frac{2k^{4}}{3} \left\{ \sigma^{(4)}(T_{2}) - \sigma^{(4)}(T_{1}) \right\},$$
(2)
$$\Delta \Phi(k,T_{1},T_{2}) \Box 2k \left\{ \sigma^{(1)}(T_{2}) - \sigma^{(1)}(T_{1}) \right\} - 4k \left\{ \frac{1}{r_{0}} + \frac{1}{\lambda} \right\} \left\{ \sigma^{2}(T_{2}) - \sigma^{2}(T_{1}) \right\} - \frac{4k^{3}}{3} \left\{ \sigma^{(3)}(T_{2}) - \sigma^{(3)}(T_{1}) \right\}.$$
(3)

where $\sigma^{(n)}$ are *n* th-order cumulants and can be expressed in terms of the power moments of the true RD function $\rho(T, r)$.

2.2. Thermodynamic parameters and anharmonic effective potential of c-Si

To determine the thermodynamic parameters of a system, it is necessary to specify its AE potential and force constants [10]. One considers a monatomic system with an anharmonic effective potential (ignored the constant contribution) is extended up to the fourth-order:

$$V_{eff}(x) = \frac{1}{2}k_0x^2 - k_3x^3 + k_4x^4, \qquad (4)$$

where $x = r - r_0$ is the deviation of the interatomic distance from the potential minimum position, k_0 is the effective force constant, k_3 and k_4 are local force constants giving asymmetry of potential due to the inclusion of anharmonicity.

The Morse potential is assumed to describe the interatomic interaction model for the potential energy of a diatomic molecule. Applying the Morse potential to calculate the interaction energy between each pair of atoms in cubic metals was proposed by Girifalco and Weizer [11]. In the present study, we expand the Morse potential to the fourth-order:

$$V(x) = D\left(e^{-2\alpha x} - 2e^{-\alpha x}\right) \cong -D + D\alpha^2 x^2 - D\alpha^3 x^3 + \frac{7}{12}D\alpha^4 x^4.$$
 (5)

where *x* is the same previously defined value, α describes the width of the potential, and *D* is the dissociation energy.

In the relative vibrations of absorbing (A) and

backscattering (B) atoms, including correlation effects and taking into account only the nearestneighbor interactions, the AE potential [6] is given by

$$V_{eff} = V(x) + \sum_{i=A,B} \sum_{j \neq A,B} V\left(\frac{\mu}{M_i} x \hat{R}_{AB} \hat{R}_{ij}\right), \tag{6}$$

where $\mu = M_A M_B / (M_A + M_B)$ is the reduced mass of the absorber and backscatterer with masses M_A and M_B , respectively, \hat{R} is a unit vector, the sum *i* is the over absorbers (i = A) and backscatterers (i = B), and the sum *j* is over the nearest neighbors.



Fig. 1. Structural model of c-Si.

Applying Eqs. (5) and (6) to the structure of c-Si with a mass of atoms is $M_A = M_B = m$, as seen in Fig. 1, the AE potential is written as

$$V_{eff}(x) = \frac{7}{3} D\alpha^2 x^2 - \frac{35}{36} D\alpha^3 x^3 + \frac{1519}{2592} \alpha^4 x^4.$$
(7)

The local force constants k_0 , k_3 , and k_4 are deduced from Eq. (7) as follows:

Comparing Eq. (4) with Eq. (7), we deduce the local force constants k_0 , k_3 , and k_4 as follows:

$$k_0 = \frac{7}{3}D\alpha^2, \quad k_3 = \frac{35}{36}D\alpha^3, \quad k_4 = \frac{1519}{2592}D\alpha^4.$$
(8)

The thermal vibration of atoms is characterized by the correlated Einstein frequency ω_E and temperature θ_E , which are calculated from the effective force constant k_0 in the following forms:

$$\omega_E = \sqrt{\frac{k_{eff}}{\mu}} = \alpha \sqrt{\frac{10D}{m}}, \quad \theta_E = \frac{h\omega_E}{k_B} = \frac{h\alpha}{k_B} \sqrt{\frac{10D}{m}}, \quad (9)$$

where k_B is the Boltzmann constant, \hbar is the reduced Planck constant.

Consequently, the correlated Einstein frequency $\omega_{\rm E}$ and temperature $\theta_{\rm E}$, and force constants

 k_0, k_3 , and k_4 are expressed in terms of the Morse potential parameters via Eqs. (8) and (9).

3. Temperature dependence of EXAFS oscillation within the classical ACE model

The EXAFS cumulants are explicitly related to low-order moments of the distribution function.

The analysis of the EXAFS spectra usually uses the first four EXAFS cumulants [10]. The expressions of the first four cumulants of diamond crystals within the classical ACE model are given as follows [8]:

$$\sigma^{(1)}; \quad \frac{15k_B}{28D\alpha}T = \frac{5\alpha}{4}\sigma^2, \tag{10}$$

$$\sigma^{(2)} \square \frac{3k_B}{7D\alpha^2} T = \frac{2k_B}{m\omega_E^2} T \equiv \sigma^2, \tag{11}$$

$$\sigma^{(3)}; \quad \frac{45k_B^2}{98D^2\alpha^3}T^2 = \frac{5\alpha}{2}(\sigma^2)^2, \tag{12}$$

$$\sigma^{(4)}; \ \frac{687k_B^3}{686D^2\alpha^4}T^3 = \frac{229\alpha^2}{18}\left(\sigma^2\right)^3.$$
(13)

Substituting the anharmonic EXAFS cumulants from Eqs. (10)-(13) into Eqs. (2) and (3) to calculated the logarithm of amplitude ratio $M(k,T_1,T_2)$ and the linear phase difference $\Delta\Phi(k,T_1,T_2)$, the results are obtained as

$$M(k,T_1,T_2) \Box -\frac{6k_B}{7D\alpha^2}k^2(T_2-T_1) + \frac{229k_B^3}{343D^3\alpha^4}k^4(T_2^3-T_1^3),$$
(14)

$$\Delta\Phi(k,T_1,T_2) \Box \frac{15k_B}{14D\alpha} k(T_2-T_1) - \frac{12k_B}{7D\alpha^2} \left\{ \frac{1}{r_0} + \frac{1}{\lambda} \right\} k(T_2-T_1) - \frac{30k_B^2}{49D^2\alpha^3} k^3 (T_2^2 - T_1^2).$$
(15)

Thus, the EXAFS cumulants using the CACE model can be expressed in the explicit and simple forms of the mean-square relative displacement (MSRD) σ^2 or temperature T. Consequently, the logarithm of amplitude ratio and the phase difference of the EXAFS oscillation is also expressed in simple forms of temperature via the EXAFS cumulants. These obtained results can describe the influence of anharmonic effects on the classical limit at high temperatures, and show that the cumulants are very useful for the quantitative treatment of the anharmonic EXAFS oscillation.

4. Numerical results and discussions

To discuss the effectiveness of the classical ACE model for the analysis of the EXAFS oscillation of c-Si in this work, we apply Eqs. (8)-

(9) in Sec. 2 and Eqs. (10)-(15) in Sec. 3 to the numerical calculations. Firstly, we perform numerical calculations for the force constants of the AE potential, the thermodynamic parameters, and the temperature dependence of the first four EXAFS cumulants in the range from 0 K to 1000 K. Our results of the cumulants are compared with those obtained using the quantum ACE model [6] in the range from 0 K to 1000 K by Hung et al. [12] and experiment at 80 K, 300 K, and 500 K by Benfatto et al. [13]. Then, we analyze the logarithm of amplitude ratio and the phase difference of the anharmonic EXAFS oscillation with reference value at θ_{E} in the wavenumber range from 0 Å⁻¹ to 20 Å-1 at 700 K, 800 K, and 900K. Lastly, we evaluate and comment on the results obtained using the classical ACE model in this work.



Fig. 2. Temperature dependence of (a) the first, (b) second, (c) third, and (d) fourth EXAFS cumulants of c-Si obtained using the ACE model, the quantum ACE model [12], and the experiments [13].

The force constants k_0 , k_3 , and k_3 , and the correlated Einstein frequency ω_E and temperature θ_E are calculated via Eqs. (8)-(9). Our obtained results are $k_0 \square 10.39 \text{ eV}\text{Å}^{-2}$, $k_3 \square 6.75 \text{ eV}\text{Å}^{-2}$, $k_4 \square 6.35 \text{ eV}\text{Å}^{-2}$, $\omega_E \square 8.42 \times 10^{13}$ Hz,

and $\theta_E \square 643.49$ K, in which the use of the Morse potential parameters D = 1.83 eV and $\alpha = 1.56$ Å⁻¹ in calculations were determined by Swalin (1961) [14].

The temperature dependence of (a) the first cumulant $\sigma^{(1)}(T)$, (b) the second cumulant



 $\sigma^2(T)$, (c) the third cumulant $\sigma^{(3)}(T)$, and (d) the fourth cumulant $\sigma^{(4)}(T)$ of c-Si is calculated by Eqs. (9)-(12) and shown in Fig. 2. Our obtained results using the classical ACE model agree well with the results obtained using the quantum ACE model [12] (for the first three cumulants) and the experiments [13] (for the second cumulants) at high temperatures, especially for high-order cumulants. For example, at 500 K, the results obtained using the classical ACE model and quantum ACE model $\sigma^{(1)} \square 8.1 \times 10^{-3} \text{ Å}, \quad \sigma^{(2)} \square 4.1 \times 10^{-3} \text{ Å}^2,$ $\sigma^{(3)} \Box 6.7 \times 10^{-5} \,\text{\AA}^3$, and $\sigma^{(4)} \Box 2, 2 \times 10^{-6} \,\text{\AA}^4$, $\sigma^{(1)} \Box 9.2 \times 10^{-3} \text{\AA}$ and $\sigma^{(2)}$ \Box 4.7×10⁻³ Å², and $\sigma^{(3)}$ \Box 6.8×10⁻⁵ Å³ [12], respectively, while the experimental value is $\sigma^{(2)} \approx 4.3 \times 10^{-3} \,\text{\AA}^2$ at 500 K [13]. Our expressions of the first three EXAFS cumulants are the same as the corresponding expressions

calculated by the quantum ACE model [12] in the

high-temperature limit. Moreover, our results of the EXAFS cumulants can agree with their experimental values at high temperatures although we have only experimental values of the second cumulant in comparisons. It is because the experimental values of other cumulants can all be deduced from the second cumulant [15].

Thus, the results of the temperature dependence of the first four EXAFS cumulants of c-Si are



Fig. 3. The (a) logarithm of amplitude ratio and (b) phase difference with reference value at θ_E of c-Si obtained using the classical ACE model at various temperatures.

The (a) logarithm of amplitude ratio $M(k,T) = \ln \left[A(k,T)/A(k,\theta_E)\right]$ and (b) phase difference $\Delta \Phi(k,T) = \Phi(k,T) - \Phi(k,\theta_E)$ of c-Si at 700 K, 800 K, and 900 K are calculated by Eqs. (14)-(15) and shown in Fig. 3. It can be seen that the values of M(k,T) and $\Delta \Phi(k,T)$ decrease with increasing temperature T and decrease with fast-increasing wavenumber k. For example, at T =700 K, T = 800 K, and T = 900 K, the approximate results of M(k,T) are -0.086, -0.231, and -0.372 at k = 10 Å⁻¹, and -0.233, -0.578, and -0.833 at k = 20 Å⁻¹, respectively, while the corresponding results of $\Delta \Phi(k,T)$ are – 0.014, – 0.043, and -0.080 at $k = 10 \text{ Å}^{-1}$, and -0.192, -

obtained using the classical ACE model, which satisfied all of their fundamental properties in comparison with the quantum ACE model and experiment at temperatures above the correlated Einstein temperature. It is explained because anharmonicity in EXAFS spectra appears from about room temperature. These results described the influence of anharmonic effects on the classical limit via thermal vibration-contributions at high temperatures.



0.573, and -1.012 at k = 20 Å⁻¹, respectively. Moreover, in the temperature dependence, the value of $\Delta \Phi(k,T)$ decreases faster than the value of M(k,T). It is because high-order cumulants increase with temperature *T*, in which the third cumulant reduces the value of $\Delta \Phi(k,T)$ and the fourth cumulant increases the value of M(k,T), as seen Eqs. (14) and (15).

Thus, the results of the temperature dependence of anharmonic EXAFS oscillation of c-Si are obtained using the classical ACE model, which shows that the even-order cumulants contribute to the amplitude reduction, and the odd-order cumulants contribute primarily to the phase shift of the anharmonic EXAFS oscillation. Accurate calculation of the cumulants will allow us to accurately analyze the change of the anharmonic EXAFS oscillation, and from which one will determine the structural parameters from the experimental EXAFS data.

5 Conclusions

In this work, the classical ACE model has been used successfully in the analysis of the anharmonic EXAFS oscillation of c-Si. The EXAFS analysis is performed based on evaluating the contribution of terms of the Debye-Waller factor in the cumulant expansion approach up to the fourth-order. The results of the first four EXAFS cumulants are not only expressed in explicit and simple forms of the temperature T or MSRD but also satisfy all of their fundamental properties in temperature dependence. The analytical results show the role and meaning of the EXAFS cumulants for the amplitude reduction and the phase shift of the anharmonic EXAFS oscillation. The obtained results are very useful for analyzing the experimental data of the anharmonic EXAFS oscillation of c-Si.

The good agreement of our numerical results with those obtained using the quantum ACD model and the experiments at various temperatures show the effectiveness of the classical ACE model for calculating and analyzing the anharmonic EXAFS oscillation of c-Si. The obtained results can be applied to the analysis of the anharmonic EXAFS oscillation of c-Si starting from about the correlated Einstein temperature to just before the melting point.

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SỰ PHỤ THUỘC VÀO NHIỆT ĐỘ CỦA DAO ĐỘNG EXAFS PHÍ ĐIỀU HÒA CỦA TINH THỂ SILIC

Tống Sỹ Tiến, Lê Việt Hoàng

Thông tin bài viết	Tóm tắt
Ngày nhận bài:	Trong công việc này, dao động của phổ cấu trúc tinh tế mở rộng (EXAFS) phi điều hòa của tinh thể silic đã được biểu diễn qua các số hang của hệ số Debye-
16/9/2020 Ngày duyết đăng:	Waller bằng phương pháp khai triển cumulant đến bậc bốn. Bốn EXAFS
10/12/2020	cumulant đầu tiên đã được tính toán dựa trên mô hình Einstein tương quan phi điều hòa (ACE) cổ điển và qui trình phân tích phù hợp, trong đó các tham số
	nhiệt động được rút ra từ hàm thế hiệu dụng phi điều hòa thu được bằng cách
Từ khóa:	tiếp cận các đóng góp lần cận của lớp nguyên từ đầu tiên. Việc phân tích sự
Phân tích EXAFS; Hê số	phụ thuộc vào nhiệt độ của dao động EXAFS được thực hiện thông qua việc
Debve-Waller: Mô hình	đánh giá ảnh hưởng của các cumulant vào sự giảm biên độ và sự dịch pha. Các
Finstein trong augn phi	kết quả tính số được tìm thấy trùng hợp tốt với các kết quả thu được bằng mô
tin the set of the set	hình ACE lượng tử và thực nghiêm ở các nhiệt đô khác nhau. Các kết quả thu
aleu noa, 1 inn ine silic.	được là rất hữu ích đối trong việc phân tích các dữ liệu EXAFS thực nghiệm
	cua c-51.